Some physical System by Fourier series in weight Functionfor Electric Fieldand its Application

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Abstract

The Fourier series is a tool for mathematics, electrical and other engineering sciences .And also it's beneficial in differential equations because it can transform them into equations which are easier to solve. Many of the phenomena studied in engineering and science are periodic in nature the current and voltage in an alternating current circuit. These periodic functions can be analyzed into their constituent components (fundamentals and harmonics) by a process called Fourier analysis. In this work, Fourier-series representation of a weight function is used to highlight and clarify the controversial problem of finding the value of the function at a point of discontinuity. Several physical situations are presented to examine the consequences of this kind of representation and its impact on some widely well-known problems whose results are not clearly understood.

1. Introduction

Jean Baptiste Joseph Fourier (1768-1830) was a French mathematician, physicist and engineer, and the founder of Fourier analysis. In 1822 he made the claim, seemingly preposterous at the time, that any function of t, continuous or discontinuous, could be represented as a linear combination of functions sin nt. This was a dramatic distinction from Taylor series. While not strictly true in fact, this claim was true in spirit and it led to the modern theory of Fourier analysis with wide applications to science and engineering. Fourier series started life as a method to solve problems about the now of heat through ordinary materials. It has grown so far that if

you search our library's catalog for the keyword \Fourier" you will nod 618 entries as of this date. It is a tool in abstract analysis and electromagnetism and statistics and radio communication and : : : . People have even tried to use it to analyze the stock market. (It didn't help.) The representation of musical sounds as sums of waves of various frequencies is an audible example. Roughly speaking, a Fourier series expansion for a function is a representation of the function as sum of sin's and cosines. Expressing a musical tone as a sum of a fundamental tone and various harmonics is such a representation. So is a spectral decomposition of light waves. The main Fourier series expansions that we use in this course are stated in the next section. We shall never prove that these expansions are correct. But in the section "Validity of Fourier series" we give an elementary partial argument that, hopefully, will convince you that the expansions are indeed correct.

The problem of a conducting sphere of radius R and uniform charge Q on its surface is discussed in introductory physics textbooks [1] [2]. For such a sphere, Gauss's law gives that the electric field, at a point r, is zero inside $(r < R$) and $\frac{kQ}{R^2}$ outside $(r > R$). Usually, the question of finding the electric field on points on the surface of the sphere is not raised due to the discontinuity of the field at such points. The limit of the electric field as r→R from inside gives zero and from outside gives $\frac{kQ}{R^2}$ and therefore the electric field is ill-defined on the surface. Most textbooks as in [1] and [2], authors calculate, by direct use of Gauss law, the electric field just outside a conducting surface and obtain the correct value $\frac{kQ}{R^2}$. This value is the same as the outside limiting value of the electric field of the conducting sphere. This causes a misleading and confusion among students for the electric field on the surface and may claim that the electric field there is $\frac{kQ}{R^2}$. This is not a correct conclusion and one should admit the ambiguity of the electric field on the surface and care must be taken when dealing with such a situation. In some cases one has to assign a value for the electric field at a very small patch on the surface for the purpose of calculating the electrostatic pressure (force per unit area) on the surface. This has been done by Griffiths in his Electrodynamics book [3]. Griffiths assigned the average value between the inside and outside limits and got the value $\frac{kQ}{2R^2}$ for the electric field on the surface of the conductor, and he arrived at the correct well-known value of the electrostatic pressure on the surface [4], namely, $\sigma^2/2\varepsilon_0$ where σ is the surface charge density and ε_0 is the electric

permittivity of free pace. Other researchers [5] [6] [7] [8] [9] considered the electric field which suffers a jump at the surface and concluded that the appropriate value of the electric field at the surface is the average value between the inside and the outside limits.

In the light of the above debate and due to the ambiguity of this problem, the aim of the present paper is to give a proof that the value of the electric field at the discontinuity point on the surface is in fact the average value between the inside and the outside limits of the electric field as r→ R. Fourier series expansion will be used to obtain our result.

2. Fourier Series Representation of a Discontinuous Function

Expansion of a function by a Fourier series has been of a great advantage in physics and engineering [10]-[18] because it allows one to more easily manipulate functions that are discontinuous or difficult to represent analytically. Some investigators have been working on fast and accurate simulations for Fourier representation of discontinuous function that represent scientific problems [19]-[25]. The value of the expanded function f (x) at its point of discontinuity point is the average of the upper and lower limits of the function [26] [27].

This means that the Fourier series converges to half-way between these two limits. Mathematically, at a point of discontinuity, x_0 the Fourier series converges to the value $f(x_0)$ given by

$$F(x_0) = \frac{1}{2} \lim_{\varepsilon \to 0} [f(x_0 + \varepsilon) + f(x_0 - \varepsilon) \tag{1}$$

The Fourier series expansion of the function ( ) fx, with period 2L is conventionally written as

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left(a_n \cos \frac{n\pi x}{L} + b_n \sin \frac{n\pi x}{L} \right) \tag{2}$$

where a_0, a_n, b_n are constants called Fourier coefficients. Due to the orthogonality of the sine and cosine functions, the Fourier coefficients are readily obtained with the result. Where

$$a_0 = \frac{1}{L} \int_{-L}^{L} f(x) dx \tag{3}$$

$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) dx \tag{4}$$

$$b_n = \frac{1}{L}\int_{-L}^{L} f(x)\sin\left(\frac{n\pi x}{L}\right) dx \tag{5}$$

The factor $\frac{1}{2}$ which appears in the a_0 term in Equation (2) is included so that Equation (4) may be applied for n=0, as well as n >0

In order to demonstrate Equation (1), we consider a specific function $f(x)$ which is weight function given by

$$f(x) = \begin{cases} 0, x \in (-L,0) \\ 1, x \in (0,L) \end{cases} \tag{6}$$

Applying Equations (3)-(5), one immediately gets $a_0 = 1, a_n, = 0, b_n = \frac{2}{n\pi}\ for\ n = odd\ only$

Therefore, Equation (2) yields

$$f(x) = \frac{1}{2} + \frac{2}{\pi}\sum_{n=odd}^{\infty}\left(\frac{1}{n}\sin\frac{n\pi x}{L}\right) \tag{7}$$

Obviously, at the point of discontinuity (x=0) the series in Equation (7)

converges to 1/2, which is the average value between the two limits of $f(x)$ below and above the jump.

3. A Uniformly Charged Conducting Sphere

In order to discuss the problem of the value of the electric field on the surface of a

conducting sphere, we consider a conducting sphere of radius R which is uniformly charged by a charge Q. Direct application of Gauss's law gives the electric field at a point r from the sphere's center, with the result

$$\breve{E} = \begin{cases} 0, \ r < R \\ \frac{kQ}{r^2}\hat{r}, \ r > R \end{cases} \tag{8}$$

with $K = \frac{1}{4\pi\varepsilon_0}$. Obviously, the electric field suffers a discontinuity at points on the surface since

its limit from inside gives zero while from outside gives $\frac{kQ}{2R^2}$

In order to find the electric field at points on the surface, we proceed as follows: We consider a conducting object which carries a surface charge density σ on its surface. The widely well-known quantity which is usually invoked is the electric field just outside the conductor. The usual treatment to derive this is by considering a cylindrical Gaussian surface whose half of its length inside the conductor and the other half is outside as in Figure 2. Application of Gauss's law yields $E = \frac{\sigma}{\varepsilon_0}$; which is found in most standard introductory physics textbooks [2]. Therefore, the electric field is zero inside the conductor and a constant ($\frac{\sigma}{\varepsilon_0}$) just outside. The result shows that the electric field suffers a finite jump on the surface, so that the electric field is written as

$$E = \begin{cases} 0, -\varepsilon \leq x < 0 \\ \dfrac{\sigma}{\varepsilon_0}, \quad 0 < x \leq \varepsilon \end{cases}$$

where ε is a small quantity and x=0 , the discontinuity point, represents a point on the surface. The point now is to expand the electric field in Fourier series over the interval [-ε ε,] , Comparing Equation (9) with Equation (6) and noting that

L $\rightarrow$ ε , the Fourier coefficients, given in Equations (3)-(5) are easily calculated

with the result $a_0 = \frac{\sigma}{\varepsilon_0}, a_n = 0, b_n = \frac{2\sigma}{n\pi\varepsilon_0}$; Therefore, Equation (2) gives the Fourier expansion of the electric field which is given by

$$E = \frac{\sigma}{2\varepsilon_0} + \frac{2\sigma}{\pi\varepsilon_0}; \sum_{n=odd}^{\infty} \left(\frac{1}{n}\sin\frac{n\pi x}{\varepsilon}\right) \tag{10}$$

It should be clear that the above series converges to $\frac{\sigma}{2\varepsilon_0}$ at the discontinuity point (x=0) which is the average value of the electric field between its two limits from inside and outside. Therefore, by using the relation $\sigma = \frac{q}{4\pi R^2}$ for the conducting sphere problem, we get our required result for the electric field on the surface,

$$E = \frac{1}{2}\left[\lim_{r} \to R \quad E\right] = \frac{Q}{2(4\pi\varepsilon_0 R)} = \frac{KQ}{2R^2} \tag{11}$$

Our result in Equation (11) removes the ambiguity of the electric field at points on the surface of a conducting sphere and assigns a value of this field at its discontinuity point. This interesting result must be explained at the undergraduate level for physics and engineering students, since it has been avoided in almost all undergraduate physics textbooks. It should be emphasized that our result, beside its mathematical interest, it also has applications in physical situations in which the value of a function at its point of discontinuity is necessary in order to derive some relevant physical quantities. In the next section, three physical situations will be presented to demonstrate the use of our result and to show that the value of a discontinuous function at its point of discontinuity is the average value of the function at that point, which is the value where the Fourier series of the function converges.

4. Some Physical Problems Involve Discontinuous Functions

Here, we consider three physical situations that involve functions that suffer a jump at a point on a boundary. Therefore, assigning a value of the function at that point is necessary in order to achieve a relevant physical quantity.

4.1. The First Problem: Electrostatic Pressure on Surface of a Conductor

The first problem deals with the calculation of electrostatic pressure on the surface of a conductor which contains a surface charge density σ. This problem has been discussed by Griffiths [28] in his famous textbook on electrodynamics. There, Griffiths considered a patch on the surface of the conductor in attempt to find the electrostatic pressure on the surface of the conductor. He argued (but not rigorously) that the electric field on the patch is the average value between the value of electric field just outside the conductor $(\frac{\sigma}{\varepsilon_0})$ and its zero value inside. Hence he arrived at the value $\frac{\sigma}{2\varepsilon_0}$ at the patch, and therefore the pressure (force per unit area) is just $\frac{\sigma^2}{2\varepsilon_0}$ However, this argument seems a bit dodgy because the patch is not a point so that part of the patch creates a field that affects the other part of the patch. So our result can now be applied to determine the value of the electric field at a point on the surface. For that purpose, we construct a small cylindrical surface as shown in Figure 1. We assume ε to be very small, so that,

$\lim\limits_{\varepsilon} 0^- \ E = 0 \ and \ \lim\limits_{\varepsilon} 0^+ \ E = \frac{\sigma}{\varepsilon_0}$ Therefore, the average of these two limits gives the   correct

value of the electric field at a point on the surface of the conducting sphere, namely, $E = \frac{\sigma}{2\varepsilon_0}$ and

hence the well-known value of the electrostatic pressure on the surface of the sphere.

4.2. The Second Problem: Energy Aspects in Charging a Capacitor

The well-known two capacitor problem has been of great interest since long time   ago, and
variety of approaches have been considered to explain the energy loss in this problem [29] [30]
[31] [32]. The essence of the problem amounts to the   problem of charging a capacitor of
capacitance C by a power supply of electromotive force V_0 with a series resistor of resistance R.
At the end of charging, the energy stored in the capacitor will be $\frac{1}{2}CV_0{}^2$ and exactly the same
amount will be dissipated regardless of the value of R. In ref. [30], the authors   used
superconducting wires and used the flux of energy carried by the Poynting vector to calculate the
energy stored in the capacitor and the energy loss. In such situation, the charge on the capacitor
involves a step function, namely, $Q(t) = Q\theta(t)$ In their derivation of their final result for the
stored energy, U, they encountered the integral

$$U = \frac{-Q^2}{C} \int \theta(t)\delta(t)dt = \frac{Q^2}{C}\theta(0) \qquad (12)$$

and a similar one for the energy loss. The value of the step function at the discontinuity point (t=
0) was used by taking the average value between the limits from below and above $t = 0$ as

$$\theta(0) = \frac{1}{2}[\theta(0^-) + \theta(0^+)] \qquad (13)$$

Therefore, the average value at the discontinuity point has a crucial role in deriving the energy
loss in this process.

4.3. The Fermi Distribution Function

In this subsection, we present our third physical situation for the use of the average value at the
jump of a discontinuous function. The average number of   fermions, in a single-particle state
with energy $\varepsilon\varepsilon$ is given by the Fermi-Dirac distribution function [33]

,

$$\langle n(\varepsilon)\rangle = \frac{1}{e^{(\varepsilon-\mu)/k_BT}+1} \qquad (14)$$

Fermi-Dirac distribution and the Fermi-level

Density of states tells us how many states exist at a given energy E. The Fermi function $f(E)$ specifies how many of the existing states at the energy E will be filled with electrons. The function $f(E)$ specifies, under equilibrium conditions, the probability that an available state at an energy E will be occupied by an electron. It is a probability distribution function.

$$f(E) = \frac{1}{1 + e^{(E-E_F)/kT}}$$

(2.7)

E_F = Fermi energy or Fermi level

k = Boltzmann constant = 1.38×10^{-23} J/K
= 8.6×10^{-5} eV/K
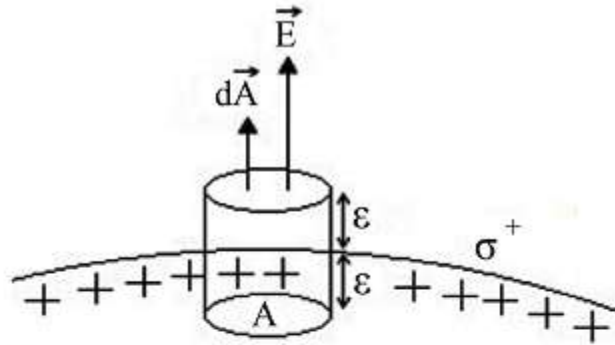
T = absolute temperature in K

1

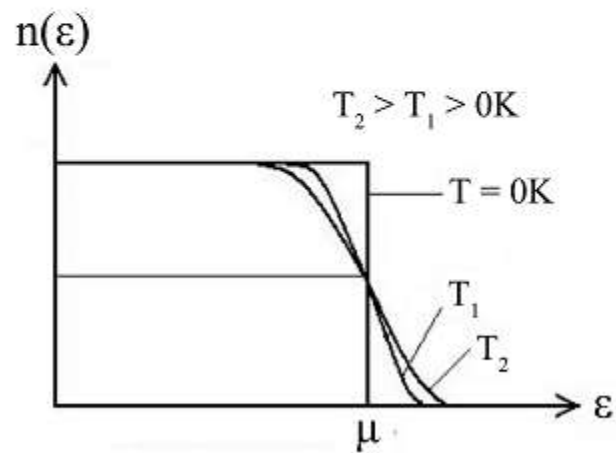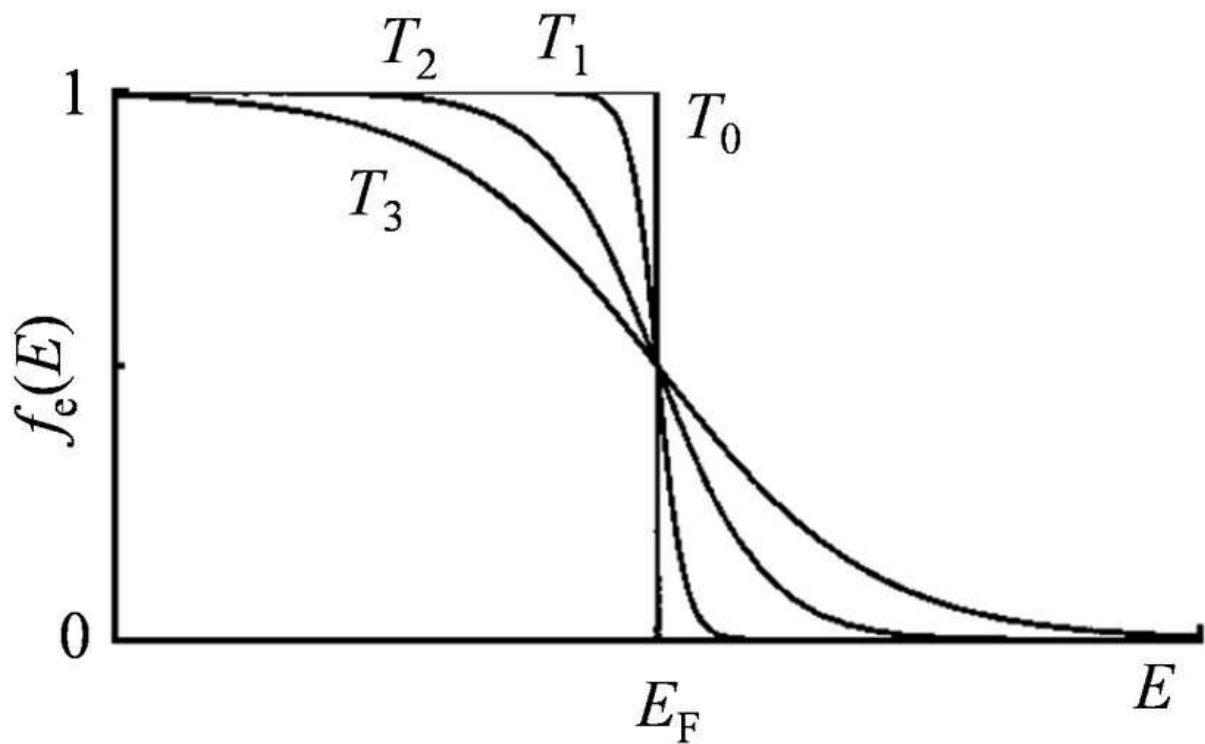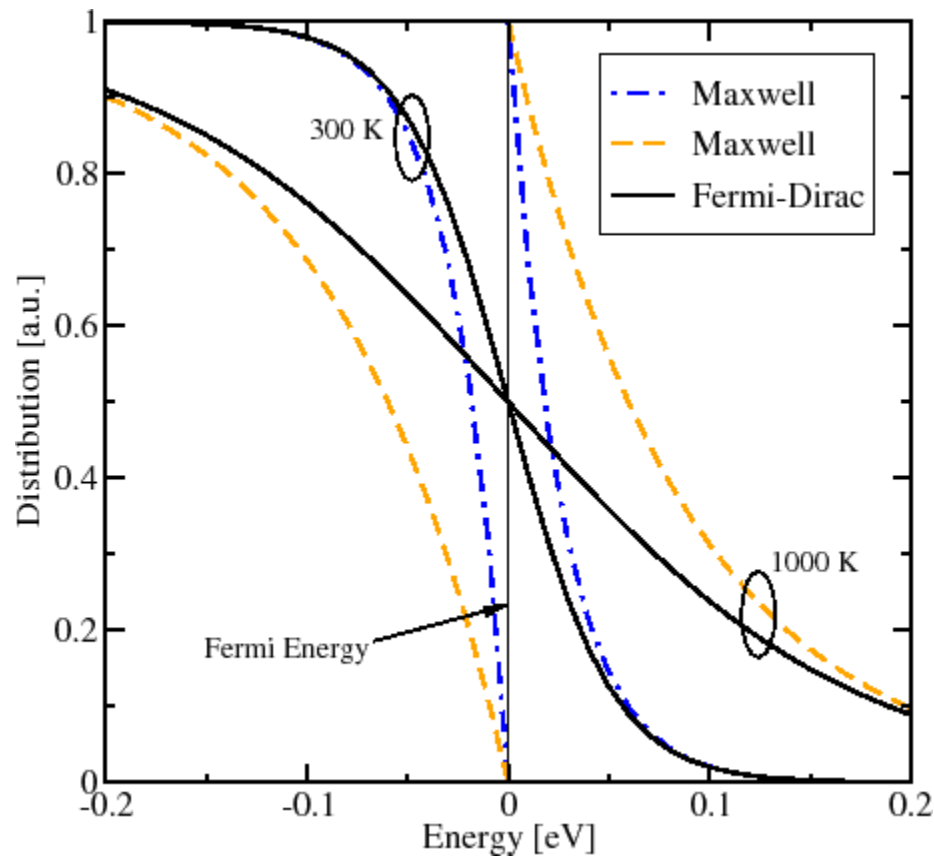**Figure 1.** Cylindrical surface on the surface of a conducting sphere.

Figure 2. Fermi distribution function as function of energy at different temperatures.

where μ is the chemical potential, B_k is the Boltzmann constant and T is the temperature in Kelvin, as shown in Figure 2. We note that the Fermi distribution function behaves like a step function

$$\lim_{T \to 0} \langle n(\varepsilon) \rangle = \begin{cases} 0 & \varepsilon > \mu \\ 1 & \varepsilon < \mu \end{cases} = \theta(\varepsilon - \mu) \qquad (15)$$

Therefore, the average population number of fermions at zero temperature is

1/2, which is at the midpoint of the jump. One can also observe form Equation (14) that for any temperature $T > 0$, the average population number is 1/2   when the energy is equal to the chemical potential.

5. Conclusion

In this paper, the authors examined the behavior of a weight function and its application in some physical systems. Fourier series representation of such function has been studied, and it has been pointed out that, at the point of discontinuity, this series converges to the average value between the two limits of the function about the jump point. So for a step function, this convergence occurs at the exact value of one half. The obtained result clarifies and solves a controversial problem about the value of the electrostatic field at points on the surface of a conducting sphere, which is usually avoided in introductory physics books. As an application of our result, three physical systems have been discussed and the average value of the function at the discontinuity point has been used in such systems: The first deals with the calculation of the electrostatic pressure on the surface of a conductor, the second concerns the calculation of the energy loss in charging a capacitor using superconducting wires and the third is the behavior of the population function of Fermi gas at zero temperature.

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